

Quantum Computing

Exercises 1: Intro to Quantum Physics

1. a) Show that the 'in' and 'out' states defined as:

$$|i\rangle = \frac{1}{\sqrt{2}}(|u\rangle + i|d\rangle)$$
$$|o\rangle = \frac{1}{\sqrt{2}}(|u\rangle - i|d\rangle)$$

are orthogonal.

b) Calculate the expectation values of σ_y in the states $|u\rangle$ and $|i\rangle$, and of σ_z in the state $|o\rangle$.

a) We take their product, which in the bracket notation reads as $\langle o|i\rangle$, and verify that it is 0:

$$\langle o|i\rangle = \frac{1}{\sqrt{2}}(\langle u| + i\langle d|) \cdot \frac{1}{\sqrt{2}}(|u\rangle + i|d\rangle) = \frac{1}{2}(\langle u|u\rangle + i\langle u|d\rangle + i\langle d|u\rangle - \langle d|d\rangle)$$

Now, recalling $|u\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$, $|d\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$, we have:

$$\langle u|u\rangle = (1 \ 0) \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 1 \quad , \quad \langle u|d\rangle = (1 \ 0) \cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix} = 0$$

$$\langle d|u\rangle = (0 \ 1) \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 0 \quad , \quad \langle d|d\rangle = (0 \ 1) \cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix} = 1$$

b) We insert the Pauli matrices into the expression for the expectation value of a general operator: $\langle \psi|A|\psi\rangle$

$$\langle u|\sigma_y|u\rangle = (1 \ 0) \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = (1 \ 0) \begin{pmatrix} 0 \\ i \end{pmatrix} = 0$$

$$\langle l|\sigma_y|l\rangle = \frac{1}{2}(1 \ -i) \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = \frac{1}{2}(1 \ -i) \begin{pmatrix} 1 \\ i \end{pmatrix} = 1$$

Alternatively, rather than working with matrix multiplication you can decompose the state into the eigenstates of the desired operator. For example

$$\langle o|\sigma_z|o\rangle = \frac{1}{2}(\langle u| + i\langle d|)\sigma_z(|u\rangle - i|d\rangle)$$

Since we know that $\sigma_z|u\rangle = |u\rangle$ and $\sigma_z|d\rangle = -|d\rangle$ we have

$$\langle o|\sigma_z|o\rangle = \frac{1}{2}(\langle u| + i\langle d|)(|u\rangle + i|d\rangle) = \frac{1}{2}(1 - 1) = 0$$

2. a) Normalise the state

$$|\psi\rangle = 3i|u\rangle + (1 - 2i)|d\rangle.$$

b) For this (normalised) state, calculate the probability of getting both positive (+1) and negative (-1) spin eigenvalues by measuring σ_z .

a) Normalization means that taking the norm of the state $|\psi\rangle$, is unity i.e. $\sqrt{\langle \psi|\psi\rangle} = 1$. In this case however, we have:

$$\langle \psi|\psi\rangle = (-3i)(3i) + (1 + 2i)(1 - 2i) = 9 + 5 = 14 \neq 1.$$

We should then rescale our state by a constant N , $|\psi\rangle \Rightarrow N \cdot |\psi\rangle$ such that the result is 1:

$$\langle N\psi|N\psi\rangle = N^2 \overbrace{\langle \psi|\psi\rangle}^{=14} = 1 \Rightarrow N = \frac{1}{\sqrt{14}}$$

So our new normalised state is

$$|\psi\rangle = \frac{3i}{\sqrt{14}}|u\rangle + \frac{(1 - 2i)}{\sqrt{14}}|d\rangle.$$

b) $P_\psi(+)$ = $|\langle u|\psi\rangle|^2 = |\frac{3i}{\sqrt{14}}|^2 = \frac{9}{14}$. For $P_\psi(-)$, we know that this is the only other possible outcome so $P_\psi(-) = 1 - \frac{9}{14} = \frac{5}{14}$.

3. (Nielsen & Chuang Ex. 2.11 [Eigendecomposition of a Pauli matrix])
 Find the eigenvectors, eigenvalues and diagonal representations of σ_x .

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}; \det(\sigma_x - \lambda I) = \begin{vmatrix} -\lambda & 1 \\ 1 & -\lambda \end{vmatrix} = \lambda^2 - 1 = 0 \Rightarrow \lambda = \pm 1$$

For the eigenvalue -1 we have:

$$\begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \cdot \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \Rightarrow \begin{cases} a + b = 0 \\ a + b = 0 \end{cases}$$

The eigenvector is therefore $v_- = \begin{pmatrix} a \\ -a \end{pmatrix}$. Choosing $a = 1$, we have $v_- = \begin{pmatrix} 1 \\ -1 \end{pmatrix}$.

And in an analogous way for the positive eigenvalue $+1$:

$$v_+ = \begin{pmatrix} 1 \\ 1 \end{pmatrix}.$$

The diagonal representation is given by:

$$\sum_i \lambda_i |v_i\rangle \langle v_i|$$

In this case we have:

$$\sigma_x = |l\rangle \langle l| - |r\rangle \langle r|$$

That is, the matrix representation of the operator in the basis v_{\pm} is the diagonal matrix, $D = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$. The change of basis matrix, P , that gets us from one matrix representation to the other has as columns the eigenvectors, $P = \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$. This is achieved in the following way:

$$P^{-1} \sigma_x P = D.$$

4. (Hermitian operators)

For a hermitian matrix \mathbf{A} , that is, a matrix that satisfies $\mathbf{A} = \mathbf{A}^\dagger$, show that:

- Different eigenvalues have orthogonal eigenvectors.
- All its eigenvalues are real. Does the converse also hold, that is, if the spectrum (the set of all eigenvalues) of a matrix is in \mathbb{R} , is it then a hermitian matrix?

a) [Done in Susskind chapter 3]

Consider two eigenvectors of A

$$A|\psi\rangle = \lambda_\psi |\psi\rangle, A|\phi\rangle = \lambda_\phi |\phi\rangle$$

Now consider:

$$\langle \psi | A | \phi \rangle = \lambda_\phi \langle \psi | \phi \rangle$$

But also since $A = A^\dagger$ this can be written as:

$$\langle \psi | A^\dagger | \phi \rangle = \lambda_\psi^* \langle \psi | \phi \rangle = \lambda_\psi \langle \psi | \phi \rangle$$

Since eigenvalues of Hermitian matrices are real as we will show below. These two expressions are equivalent so subtracting one from the other leads to:

$$(\lambda_\phi - \lambda_\psi) \langle \psi | \phi \rangle = 0$$

$\lambda_\psi \neq \lambda_\phi$ so their difference is $\neq 0$. Therefore the eigenvectors must be orthogonal for the above to be true.

b) Consider the matrix element of the adjoint of the operator: $\langle \phi | A^\dagger | \psi \rangle$

The operator can either act on the ket (that is, from the left) or on the bra, in which case it is 'daggered':

$$\langle \phi | A^\dagger | \psi \rangle = \langle A \phi | \psi \rangle$$

Now, we know that for any bra(c)ket we have: $\langle \phi | \psi \rangle = \langle \psi | \phi \rangle^*$

Taking this remark into account:

$$\langle \phi | A^\dagger | \psi \rangle = \langle A \phi | \psi \rangle = \langle \psi | A \phi \rangle^* = \langle \psi | A | \phi \rangle^*$$

Particularising for the case where $\phi = \psi$ and taking into account the eigenvalue equation, $\mathbf{A}|\psi\rangle = a|\psi\rangle$, we retrieve the eigenvalues of the operator:

$$\langle\psi|\mathbf{A}|\psi\rangle = \langle\psi|a|\psi\rangle = a \cdot \langle\psi|\psi\rangle = a$$

Last, by assumption, we have $A^\dagger = A$, so:

$$\langle\psi|A^\dagger|\psi\rangle = \langle\psi|A|\psi\rangle = \langle\psi|A|\psi\rangle^* \Rightarrow a = a^*$$

The converse is not true in general, e.g. $\begin{pmatrix} 2 & 1 \\ 0 & 1 \end{pmatrix}$

5. (Unitary operators)

Now, consider a unitary matrix, one for which

$$UU^\dagger = \mathbb{I} \iff U^\dagger U = \mathbb{I} \iff U^{-1} = U^\dagger$$

holds. Prove that its eigenvalues are of the form $e^{i\theta}$ and that eigenvectors of different eigenvalues must be orthogonal as well.

We start by writing the eigenvalue equation for the unitary operator U :

$$U|x\rangle = \lambda_x|x\rangle \leftrightarrow \langle x|U^\dagger = \langle x|\lambda_x^*$$

We have normalised our eigenvector, so:

$$\langle x|x\rangle = 1 = \langle x|U^\dagger U|x\rangle = \langle x|\lambda_x^* \lambda_x|x\rangle = |\lambda_x|^2 \langle x|x\rangle = |\lambda_x|^2 \Rightarrow |\lambda_x|^2 = 1 \Rightarrow \lambda_x = e^{i\theta}$$

Now, as in the previous example, consider two eigenvectors of U :

$$U|x\rangle = \lambda_x|x\rangle, \quad U|y\rangle = \lambda_y|y\rangle$$

If $\lambda_x \neq \lambda_y$, then:

$$\langle x|y\rangle = \langle x|U^\dagger U|y\rangle = \lambda_x^* \lambda_y \langle x|y\rangle \Rightarrow \langle x|y\rangle (1 - \lambda_x^* \lambda_y) = 0 \xrightarrow{\lambda_x^* \neq \lambda_y} \langle x|y\rangle (\lambda_x - \lambda_y) = 0$$